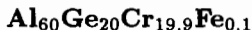


PHYSICAL PROPERTIES OF METASTABLE ICOSAHEDRAL



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The results of X-ray diffraction, ^{57}Fe Mössbauer spectroscopy, electrical resistivity, and magnetic susceptibility studies of the metastable icosahedral alloy $\text{Al}_{60}\text{Ge}_{20}\text{Cr}_{19.9}\text{Fe}_{0.1}$ are presented. The observed broadening of the diffraction Bragg peaks reflects the presence of the topological/chemical disorder. The distribution of the electric quadrupole interactions indicates the existence of a multiplicity of Fe sites. Although the low-temperature value of the electrical resistivity ($400 \mu\Omega \text{ cm}$) is rather small, we found that the temperature coefficient of the electrical resistivity is negative. The temperature dependence of the magnetic susceptibility is well described by the Curie-Weiss law with the magnetic moments localized at the transition-metal atoms.

1 Introduction

Soon after discovery of the first icosahedral (*i*) alloys in the binary Al-TM (TM=transition metal) system, it was found that the addition of a metalloid (Me) improves their structural quality.^{1,2} These ternary Al-TM-Me *i* alloys were mainly studied with respect to their structural properties. The purpose of this paper is to present the results of experimental studies of the physical properties of the $\text{Al}_{60}\text{Ge}_{20}\text{Cr}_{19.9}\text{Fe}_{0.1}$ *i* alloy with several techniques. It will be shown that this alloy, in spite of its metastable nature and a significant amount of structural disorder, exhibits a negative temperature coefficient of the electrical resistivity characteristic of high-quality stable *i* alloys.

2 Experimental

An ingot of composition $\text{Al}_{60}\text{Ge}_{20}\text{Cr}_{19.9}\text{Fe}_{0.1}$ was prepared by arc melting in an argon atmosphere of high-purity Al, Ge, Cr, and Fe enriched to 95% in

an ^{57}Fe isotope. The ingot was melt spun in air by ejecting molten alloy at 1423(10) K through a 0.7 mm orifice in a quartz tube onto a surface of a copper wheel rotating with a tangential velocity of 71(1) m/s. The resulting ribbons were about 2 cm long and 2 mm wide.

X-ray diffraction (XRD) measurements were performed on a Philips X'Pert scanning diffractometer using Cu $K\alpha$ radiation, the $K\beta$ line being eliminated by using a Kevex PSi2 Peltier cooled Si detector. The ^{57}Fe Mössbauer spectroscopy (MS) measurement at room temperature was conducted using a standard Mössbauer spectrometer operating in a sine mode. The surface density of the Mössbauer absorber was $34 \mu\text{g } ^{57}\text{Fe}/\text{cm}^2$ and it can be regarded as being thin. The electrical resistivity measurement was done with a standard dc four-probe method between 10 and 280 K. Magnetic susceptibility was measured with a SQUID magnetometer in a field of 1 kOe between 4.2 and 300 K.

3 Results and Discussion

The XRD pattern (Fig. 1) exhibits broadened Bragg peaks due to the i phase. The two most prominent i lines are on a broad background, which indicates the presence of some second phase in the studied sample.³

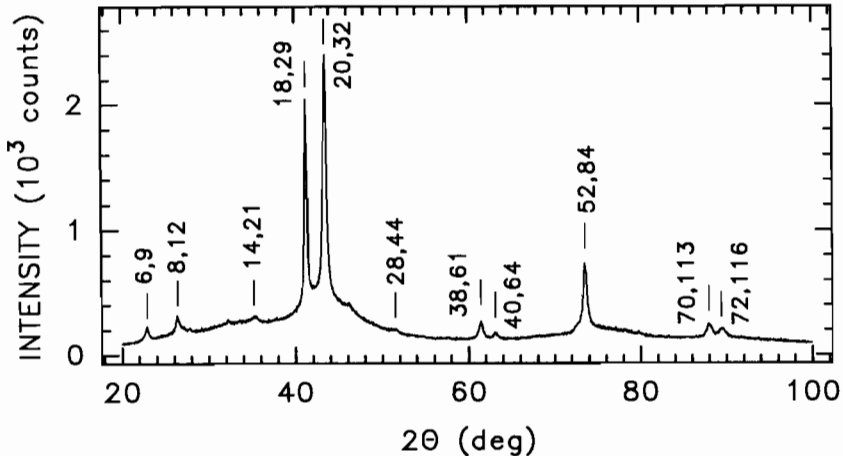


Figure 1: X-ray diffraction pattern of an $\text{Al}_{60}\text{Ge}_{20}\text{Cr}_{19.9}\text{Fe}_{0.1}$ alloy. The vertical lines above the Bragg peaks label the i peaks using the indexing scheme of Cahn *et al.* (Ref. 4).

A Mössbauer spectrum of i $\text{Al}_{60}\text{Ge}_{20}\text{Cr}_{19.9}\text{Fe}_{0.1}$ (Fig. 2) consists of two

broad and structureless lines. These broad lines result from the distribution of the electric quadrupole splittings, $P(\Delta)$.⁵ The spectrum can be fitted relatively well with a structureless Gaussian-like $P(\Delta)$ [Figs. 2(a) and (c)]. A close inspection of the residuals in Fig. 2(a) indicates the presence of some structure due to an unknown impurity component, which is consistent with the XRD result. To a first approximation, this component can be taken in the form of a symmetric Lorentzian doublet. Its inclusion in the spectral analysis [Fig. 2(b)] leads to a perfect fit. The occurrence of $P(\Delta)$ in the studied sample [Fig. 2(c)] proves that the Fe atoms are distributed among the multiplicity of sites which must result from the presence of a chemical/topological disorder in the investigated sample.

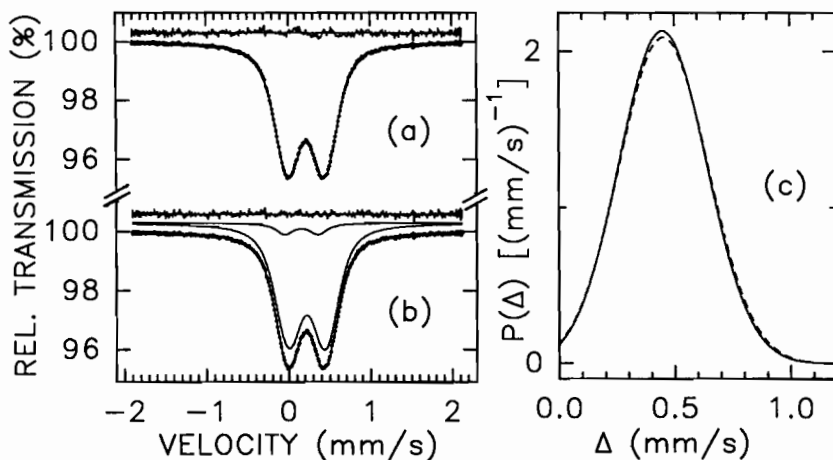


Figure 2: Room-temperature Mössbauer spectrum of $i\text{Al}_{60}\text{Ge}_{20}\text{Cr}_{19.9}\text{Fe}_{0.1}$ fitted with (a) one $P(\Delta)$ component [the solid line in (c)] and with (b) the $P(\Delta)$ [the broken line in (c)] plus an impurity component. The component spectra are also shown in (b). The residuals, multiplied by a factor of three, are shown above each spectrum. The velocity scale is relative to $\alpha\text{-Fe}$.

The electrical resistivity, ρ , of $i\text{Al}_{60}\text{Ge}_{20}\text{Cr}_{19.9}\text{Fe}_{0.1}$ (Fig. 3) is much larger than that of ordinary metallic alloys. In spite of the presence of significant structural disorder in the sample, the $\rho(T)$ dependence (Fig. 3) is similar to that observed for high-quality stable i alloys.⁶ A relatively large value of ρ and an enhanced thermal stability of i Al-TM-Me alloys^{1,2} can be explained within the framework of the Hume-Rothery mechanism.⁷ For the studied alloy the electron-per-atom ratio is 1.67. This leads to the diameter of the Fermi

sphere being 2.88 \AA^{-1} which compares well with the wave number of $2.87(1) \text{ \AA}^{-1}$ corresponding to the 18,29 *i* peak (Fig. 1).

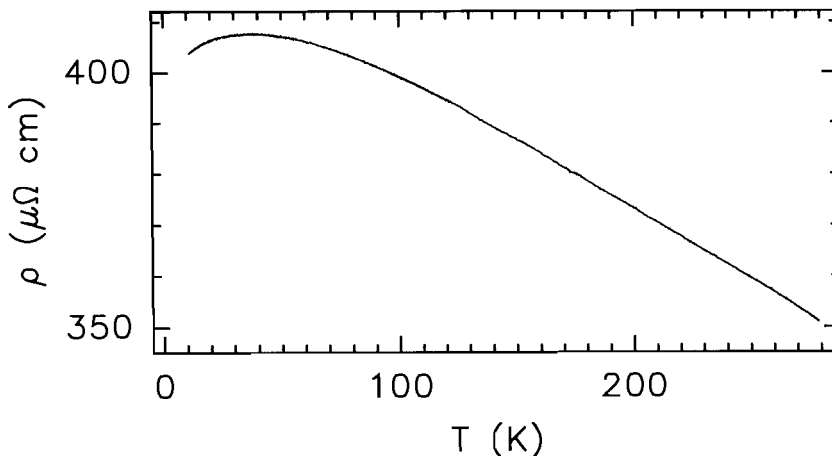


Figure 3: Temperature dependence of the electrical resistivity of *i* $\text{Al}_{60}\text{Ge}_{20}\text{Cr}_{19.9}\text{Fe}_{0.1}$.

The temperature dependence of the magnetic susceptibility, χ , follows a Curie-Weiss law $\chi = \chi_o + C/(T-\Theta)$, where the symbols χ_o , C , and Θ are respectively the temperature-independent magnetic susceptibility, the Curie constant, and the paramagnetic Curie temperature, only for temperatures larger than about 50 K (Fig. 4). The values of χ_o , C , and Θ obtained from the fit are respectively $0.370(1) \times 10^{-5} \text{ cm}^3\text{g}^{-1}$, $31.7(3) \times 10^{-5} \text{ cm}^3\text{g}^{-1}\text{K}$, and $-27.0(6) \text{ K}$. This value of C corresponds to the effective magnetic moment per TM atom of $0.708(3) \mu_B$. A deviation from the Curie-Weiss law at low temperatures (Fig. 4) may be due to the presence of some magnetic impurity in the sample, as indicated by the X-ray diffraction and Mössbauer data. The presence of such an impurity would also explain the value of χ_o , which is about an order of magnitude larger than that found in other Al-Cr-TM-Me *i* alloys.^{8,9} The negative value of Θ indicates the predominantly antiferromagnetic interaction between the TM atoms. The value of the effective magnetic moment per TM atom is similar to that observed for other Al-Cr-TM-Me *i* alloys.^{8,9} Its origin is due to the topological disorder^{9,10} present in the studied sample, as evidenced in the X-ray diffraction and Mössbauer spectra. This disorder is caused by both the quasiperiodic structure itself, in which no two atoms have exactly the same

environment, and by the rapid-quenching process involved in the production of the metastable sample.

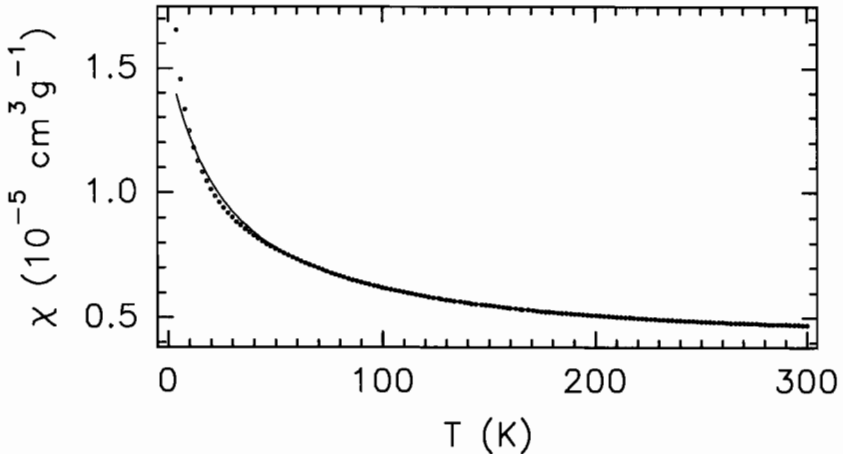


Figure 4: Temperature dependence of the magnetic susceptibility of i $\text{Al}_{60}\text{Ge}_{20}\text{Cr}_{19.9}\text{Fe}_{0.1}$. The solid line is the fit to a Curie-Weiss law for the susceptibility data at temperatures ≥ 50 K.

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