QUANTUM THEORY GROUP: FROM DISCOVERY TO INNOVATION

QNANO: COMPUTATIONAL PLATFORM FOR ELECTRONIC PROPERTIES OF SEMICONDUCTOR AND GRAPHENE NANOSTRUCTURES

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QNANO:COMPUTATIONAL PLATFORM FOR ELECTRONIC PROPERTIES OF QUANTUM NANOSTRUCTURES

OUTLINE:

INTRODUCTION ATOMIC STRUCTURE DEFINITION STRAIN ELECTRONIC STRUCTURE CALCULATION TIGHT BINDING METHOD **EFFECT OF STRAIN** SURFACE PASSIVATION **EXTERNAL FIELDS MANY-BODY EFFECTS** MULTI-EXCITON COMPLEXES **CHARGED EXCITONS EXAMPLES: InAs/GaAs SELF-ASSEMBLED DOTS**

CdSe NANOCRYSTAL

REFERENCES

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REFERENCES

QNANO SEMICONDUCTOR AND CARBON NANOSTRUCTURES

InAs/GaAs SELF-ASSEMBLED QDOTS: 2 MLN ATOMS, ~10⁷ ELE.



NANOCRYSTALS CdSe,PbS,. ~1000 ATOMS, ~10⁴ ELE. PLUS LIGAND:



InAs/InP QDOTS ON PATTERNED SUBSTRATES: 10 MLN ATOMS, ~10⁸ ELE.

GRAPHENE QDOTS: 100nm*100nm, ~10⁶ ELE.



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pyramid

ATOMISTIC THEORY OF QDOTS

NO EXACT SOLUTION TO THIS MANY BODY PROBLEM EXISTS

NO AB-INITIO APPROXIMATE SOLUTION EXISTS

APPROXIMATE METHODS:

TIGHT BINDING APPROACH

Bryant et al. Klimeck et al. Whaley et al.

• • • • • • • • • • • •

TB-DFT APPROACH Frauenheim et al.

PSEUDOPOTENTIAL APPROACH Zunger et al.







QNANO-ATOMISTIC THEORY OF QDOTS

QNANO:

- É MATERIALS:InAs,GaAs,InP,CdSe,PbSe,C,ZnTe,...)
- É BULK ELECTRONIC STRUCTURE
- É MILLION ATOM QDOTS
- É SPECIFIC CRYSTAL LATTICES
- É INTERFACES
- É STRAIN MULTIPLE SCALES
- É SPIN ORBIT COUPLING
- É ATOMISTIC DISORDER
- É EXCITON FINE STRUCTURE
- É POLARIZATION OF LIGHT
- É MULTI-EXCITON COMPLEXES
- É MAGNETIC FIELD
- É MAGNETIC IMPURITIES É



OTHER WORK: tb:BRYANT,KLIMECK,...., pseudopotential: ZUNGER,...

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REFERENCES



QNANO: BUILDING QUANTUM DOT WITH ATOMS

Zinc blend lattice



Build GaAs layer

Replace Ga with In Put InAs On GaAs Lattice



QNANO: ATOMISTIC CALCULATION OF STRAIN



MINIMIZE TOTAL ELASTIC ENERGY FIND POSITION OF ATOMS

$$E_{TOT} = E_{TOT} \left(\mathbf{R}_{1} \mathbf{R}_{2} \mathbf{R}_{3}, \dots, \mathbf{R}_{N} \right)$$

VALENCE FORCE FIELD:

$$2 \sum_{i} \sum_{j=1}^{A} \prod_{ij} \left(\left(\mathbf{R}_{i} - \mathbf{R}_{j} \right)^{2} \right)^{2}$$
$$+ \sum_{i} \sum_{j=1}^{3} \sum_{k=j+1}^{4} B_{ijk} \left(\left(\mathbf{R}_{j} - \mathbf{R}_{i} \right) \left(\mathbf{R}_{k} - \mathbf{R}_{i} \right) - \cos \theta d_{ij}^{0} d_{ik}^{0} \right)^{2}$$

 10^8 atoms = 300.000.000 variables

OPTIMAL DISPLACEMENT FIELD FOUND WITH CONJUGATED-GRADIENT METHOD (P. Keating, A. Zunger, C. Pryor,...)





M.Zielinski

STRAIN RELAXED STRUCTURE



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ELECTRONIC STRUCTURE CALCULATION TIGHT BINDING APPROACH

Quasi-electron



$$-\nabla^2/2 + V_{ion}(r) + V_{Hartree}(r) + \Sigma(r, E_l) \phi_l(r) = E_l \phi_l(r)$$

LCAO APPROACH $\varphi_l(r) = \sum_{l=1}^{N} \sum_{l=1}^{20} \phi_{l,\alpha}(R) u_{\alpha}(r-R)$ $R=1 \alpha=1$ Sp3d5s*

> We dongt really know electronic density nor V_h and V_{xc} ! Parametrize matrix elements

$$< u_b(R') | H | u_\alpha(R) >$$

Tight binding Hamiltonian



orbitals

Slater-Koster



ELECTRONIC STRUCTURE CALCULATION TIGHT BINDING APPROACH - ORBITALS



ENERGY



TIGHT-BINDING PARAMETRIZATION WITH 1st NEAREST NEIGHBORS

PROBLEM: 100 PARAMETERS!



ELECTRONIC STRUCTURE CALCULATION TIGHT BINDING MATRIX ELEMENTS

ZINCBLENDE LATTICE: TETRAHEDRAL COORDINATION



ELECTRONIC STRUCTURE CALCULATION TIGHT BINDING MATRIX ELEMENTS

TABLE I. Energy integrals for crystal in terms of two-center integrals.

	$E_{s,s}$	$(ss\sigma)$		
	$E_{s,x}$	$l(sp\sigma)$		
	$E_{x,x}$	$l^{2}(pp\sigma)+(1-l^{2})(pp\pi) \leftarrow l^{2}(pp\sigma)+(1-l^{2})(pp\pi)$		
	$E_{x,y}$	$lm(pp\sigma) - lm(pp\pi)$		
	$E_{x,z}$	$ln(pp\sigma) - ln(pp\pi)$		
	$E_{s,xy}$	$\sqrt{3}lm(sd\sigma)$		
	E_{s,x^2-y^2}	$\frac{1}{2}\sqrt{3}\left(l^2-m^2\right)\left(sd\sigma\right)$		
	Es. 322-r2	$[n^2 - \frac{1}{2}(l^2 + m^2)](sd\sigma) = 100 \text{ RULES}$		
	$E_{x,xy}$	$\sqrt{3}l^2m(pd\sigma) + m(1-2l^2)(pd\pi)$		
	$E_{x, yz}$	$\sqrt{3}lmn(pd\sigma) - 2lmn(pd\pi)$ REDUCE 100 HOPPING TERMS		
	$E_{x, zx}$	$\sqrt{3}l^2n(pd\sigma) + n(1-2l^2)(pd\pi)$		
	E_{x, x^2-y^2}	$\frac{1}{2}\sqrt{3}l(l^2 - m^2)(pd\sigma) + l(1 - l^2 + m^2)(pd\pi) $ TO 21 NONTRIVIAL PARAMETERS		
	E_{y, x^2-y^2}	$\frac{1}{2}\sqrt{3}m(l^2-m^2)(pd\sigma)-m(1+l^2-m^2)(pd\pi)$		
	E_{z,x^2-y^2}	$\frac{1}{2}\sqrt{3}n(l^2-m^2)(pd\sigma)-n(l^2-m^2)(pd\pi)$		
	$E_{x, 3z^2-r^2}$	$l[n^2 - \frac{1}{2}(l^2 + m^2)](pd\sigma) - \sqrt{3}ln^2(pd\pi)$		
	$E_{y, 3z^2 - r^2}$	$m[n^2 - \frac{1}{2}(l^2 + m^2)](pd\sigma) - \sqrt{3}mn^2(pd\pi)$		
	$E_{z, 3z^2-r^2}$	$n[n^2 - \frac{1}{2}(l^2 + m^2)](pd\sigma) + \sqrt{3}n(l^2 + m^2)(pd\pi)$		
	$E_{xy,xy}$	$3l^2m^2(dd\sigma) + (l^2 + m^2 - 4l^2m^2)(dd\pi) + (n^2 + l^2m^2)(dd\delta)$		
	$E_{xy, yz}$	$3lm^2n(dd\sigma) + ln(1-4m^2)(dd\pi) + ln(m^2-1)(dd\delta)$		
	$E_{xy, zx}$	$3l^2mn(dd\sigma) + mn(1-4l^2)(dd\pi) + mn(l^2-1)(dd\delta)$		
	E_{xy, x^2-y^2}	$\frac{3}{2}lm(l^2 - m^2)(dd\sigma) + 2lm(m^2 - l^2)(dd\pi) + \frac{1}{2}lm(l^2 - m^2)(dd\delta)$		
	E_{yz,x^2-y^2}	$\frac{3}{2}mn(l^2-m^2)(dd\sigma) - mn[1+2(l^2-m^2)](dd\pi) + mn[1+\frac{1}{2}(l^2-m^2)](dd\delta)$		
	E_{zx,x^2-y^2}	$\frac{3}{2}nl(l^2 - m^2)(dd\sigma) + nl[1 - 2(l^2 - m^2)](dd\pi) - nl[1 - \frac{1}{2}(l^2 - m^2)](dd\delta)$		
	$E_{xy, 3z^2-r^2}$	$\sqrt{3}lm[n^2 - \frac{1}{2}(l^2 + m^2)](dd\sigma) - 2\sqrt{3}lmn^2(dd\pi) + \frac{1}{2}\sqrt{3}lm(1 + n^2)(dd\delta)$		
	$E_{yz, 3z^2-r^2}$	$\sqrt{3}mn[n^2 - \frac{1}{2}(l^2 + m^2)](dd\sigma) + \sqrt{3}mn(l^2 + m^2 - n^2)(dd\pi) - \frac{1}{2}\sqrt{3}mn(l^2 + m^2)(dd\delta)$		
	$E_{zz, 3z^2-r^2} \qquad \qquad \sqrt{3}ln[n^2 - \frac{1}{2}(l^2 + m^2)](dd\sigma) + \sqrt{3}ln(l^2 + m^2 - n^2)(dd\pi) - \frac{1}{2}\sqrt{3}ln(l^2 + m^2)(dd\delta)$			
	$E_{x^2-y^2, x^2-y^2}$	$\frac{3}{4}(l^2 - m^2)^2(dd\sigma) + \left[l^2 + m^2 - (l^2 - m^2)^2\right](dd\pi) + \left[n^2 + \frac{1}{4}(l^2 - m^2)^2\right](dd\delta)$		
	$E_{x^2-y^2,3z^2-r^2}$	$\frac{1}{2}\sqrt{3}(l^2 - m^2)\left[n^2 - \frac{1}{2}(l^2 + m^2)\right](dd\sigma) + \sqrt{3}n^2(m^2 - l^2)(dd\pi) + \frac{1}{4}\sqrt{3}(1 + n^2)(l^2 - m^2)(dd\delta)$		
	$E_{3z^2-r^2, 3z^2-r^2}$	$[n^{2}-\frac{1}{2}(l^{2}+m^{2})]^{2}(dd\sigma)+3n^{2}(l^{2}+m^{2})(dd\pi)+\frac{3}{4}(l^{2}+m^{2})^{2}(dd\delta)$		

J. SLATER & G. KOSTER, PHYS. REV. 94, 1498 (1954)

ELECTRONIC STRUCTURE CALCULATION Sp3s* TIGHT BINDING HAMILTONIAN



SPIN-ORBIT INTERACTION AFTER CHADI ET AL., PHYS. REV. B



TIGHT-BINDING PARAMETERS FROM BULK BAND STRUCTURE



ELECTRONIC STRUCTURE CALCULATION TIGHT BINDING BULK BANDSTRUCTURE

TIGHT-BINDING PARAMETERS FROM BULK BAND STRUCTURE



BLOCH STATES

$$\left| \Psi_{ANION} \right\rangle = \sum_{\alpha=1}^{20} A_{\alpha}^{ANION} \left(\frac{1}{\sqrt{N}} \sum_{i, BLUE} e^{ikR_i} u_{\alpha} \left(\vec{r} - \vec{R}_i \right) \right)$$
$$\left| \Psi_{CATION} \right\rangle = \sum_{\alpha=1}^{20} B_{\alpha}^{CATION} \left(\frac{1}{\sqrt{N}} \sum_{j, RED} e^{ikR_j} u_{\alpha} \left(\vec{r} - \vec{R}_j \right) \right)$$

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ELECTRONIC STRUCTURE CALCULATION FITTING BULK BANDSTRUCTURE



ENERGIES AND AVAILABLE EFFECTIVE MASSES FITTED USING GENETIC ALGORITHM

EXAMPLE: InAs

NONLOCAL PSEUDOPOTENTIAL, CHELIKOVSKY & COHEN, PRB

ELECTRONIC STRUCTURE CALCULATION TIGHT BINDING BULK BANDSTRUCTURE





Tight binding

KORKUSINSKI, NRC1

ELECTRONIC STRUCTURE CALCULATION TIGHT BINDING BULK BANDSTRUCTURE

2		TARGET	FITTED
1- InAs	${ m E_g}_{ m SO}$	0.418 eV 0.38 eV	0.413 eV 0.381 eV
	m _e	0.024 m ₀	$0.025 \ m_0$
-1	m _{lh}	0.026 m ₀	0.029 m ₀ [100] 0.029 m ₀ [111]
-3-	m _{hh [100]} m _{hh [111]}	$0.35 m_0$ $0.43 m_0$	0.369 m ₀ 0.471 m ₀
4	m _{so}	0.14 m ₀	0.098 m ₀
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INTERFACES





ÉTB PARAMETERS ACROSS THE INTERFACE -INTERFACE PARAMETERS – NOT A PROBLEM FOR GaAs/InAs -BAND OFFSETS IN DIAGONAL MATRIX ELEMENTS

ÉMODIFIED ATOMIC POSITIONS



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REFERENCES

TB HAMILTONIAN WITH STRAIN

RESCALE MATRIX ELEMENTS BY DIRECTIONAL COSINES AND BOND STRETCHING TERMS (HARRISON'S LAW)

NO STRAIN

WITH STRAIN

$$\vec{d} = d_0 (l\vec{x} + m\vec{y} + n\vec{z})$$

$$\vec{d}' = d' (l'\vec{x} + m'\vec{y} + n'\vec{z})$$

$$t_{sa,zc} = nV_{sa,pc\sigma}$$

$$\vec{d}' = d' (l'\vec{x} + m'\vec{y} + n'\vec{z})$$

$$\vec{d}' = d' (l'\vec{x} + m'\vec{y} + n'\vec{z})$$



LOEWDIN-ORTHOGONALIZED BASIS:

$$\left|\Phi\right\rangle_{i}=\left|R,\alpha\right\rangle-\frac{1}{2}\sum_{R',\beta}S_{R\alpha,R',\beta}\left|R',\beta\right\rangle$$

S MATRIX CHANGES AS ATOMS ARE DISPLACED

T.B. BOYKIN, G. KLIMECK, PRB



DIAGONAL HAMILTONIAN ELEMENTS

$$\varepsilon_{R\alpha} = \langle \Phi | H | \Phi \rangle_{i} \approx \langle R, \alpha | H | R, \alpha \rangle - \sum_{R', \beta} S_{R\alpha, R', \beta} \langle R, \alpha | H | R', \beta \rangle$$

SENSITIVE TO STRAIN!

T.B. BOYKIN, G. KLIMECK, PRB

SHIFTS OF DIAGONAL TERMS

PROCEDURE: USE EXTENDED HUECKEL RULE:

$$S_{R\alpha,R'\beta} \approx A \frac{t_{R\alpha,R'\beta}}{\varepsilon_{R\alpha}^{ATOM} + \varepsilon_{R'\beta}^{ATOM}} \quad \tilde{o}BARE\tilde{o} \text{ ATOMIC ENERGIES}$$

$$\varepsilon_{\vec{R}\alpha} = \varepsilon_{\vec{R}\alpha}^{0} + \sum_{\vec{R}' \in nn} \sum_{\beta} C_{(\vec{R}\alpha, \vec{R}'\beta)} \frac{\left(t_{\vec{R}\alpha, \vec{R}'\beta}^{0}\right)^{2} - \left(t_{\vec{R}\alpha, \vec{R}'\beta}^{'}\right)^{2}}{\varepsilon_{\vec{R}\alpha}^{0} + \varepsilon_{\vec{R}'\beta}^{0} - \Delta E_{ATOM}}$$

C ó PARAMETERS TO BE FITTED

FITTING OF STRAIN

HYDROSTATIC Ó BIR-PIKUS MODEL



LINEAR SHIFTS OF BAND EDGES

$$E_{CB} = E_{CB}^{0} + a_{C} \left(\varepsilon_{xx} + \varepsilon_{yy} + \varepsilon_{zz} \right)$$

$$E_{VB,HH} = E_{VB,HH}^{0} + a_V \left(\varepsilon_{xx} + \varepsilon_{yy} + \varepsilon_{zz} \right)$$

$$E_{SO} = E_{SO}^{0} + a_{V} \left(\varepsilon_{xx} + \varepsilon_{yy} + \varepsilon_{zz} \right)$$

 a_C, a_V - DEFORMATION POTENTIALS



FITTING OF STRAIN




BAND EDGES AS A FUNCTION OF STRAIN





Kadantsev et al. JAP2010

STRAIN IN TIGHT BINDING MODEL

SUMMARY : STRAIN EFFECTS CHANGE BONDS ANGLES AND LENGTHS. Bonds angles: Slater-Koster formalism

Bonds lengths: a generalized version of Harrison law

$$V_{\alpha\beta\gamma} = V^0_{\alpha\beta\gamma} \left(d_{ij}^{0} / d_{ij} \right)^{\eta_{\alpha\beta\gamma}}$$

 d_{ij} is bond length, 0 superscript = values without strain

Diagonal elements vary in response to displacement of neighbours:

$$\mathcal{E}_{\vec{R}\alpha} = \mathcal{E}_{\vec{R}\alpha}^{0} + \sum_{\vec{R}' \in nn} \sum_{\beta} C_{(\vec{R}\alpha, \vec{R}'\beta)} \frac{\left(t_{\vec{R}\alpha, \vec{R}'\beta}^{0}\right)^2 - \left(t_{\vec{R}\alpha, \vec{R}'\beta}\right)^2}{\mathcal{E}_{\vec{R}\alpha}^0 - \mathcal{E}_{\vec{R}'\beta}^0}$$

diagonal and t off-diagonal Hamiltonian matrix eler , C empirical parameters





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> InAs/GaAs SELF-ASSEMBLED DOTS CdSe NANOCRYSTAL

REFERENCES



1. ROTATE THE sp³ BASIS INTO HYBRIDIZED ORBITALS $|sp_{1}^{3}\rangle = \frac{1}{2}(|s\rangle + |p_{x}\rangle + |p_{y}\rangle + |p_{z}\rangle) \qquad |sp_{2}^{3}\rangle = \frac{1}{2}(|s\rangle + |p_{x}\rangle - |p_{y}\rangle - |p_{z}\rangle)$ $|sp_{3}^{3}\rangle = \frac{1}{2}(|s\rangle - |p_{x}\rangle + |p_{y}\rangle - |p_{z}\rangle) \qquad |sp_{4}^{3}\rangle = \frac{1}{2}(|s\rangle - |p_{x}\rangle - |p_{y}\rangle + |p_{z}\rangle)$

2. IDENTIFY DANGLING BONDS, SHIFT THEIR ENERGIES

3. ROTATE HYBRIDIZED ORBITALS BACK TO sp³ BASIS

SURFACE PASSIVATION-BOUNDARIES



IS IT IMPORTANT?

LANCZOS ITERATIVE DIAGONALIZATION







$$\varepsilon_{\alpha\vec{R}}^{0} = \left\langle \alpha\vec{R} \left| \mathbf{H}^{0} \right| \alpha\vec{R} \right\rangle$$
$$t_{\alpha\vec{R},\alpha'\vec{R}'}^{0} = \left\langle \alpha'\vec{R}' \left| \mathbf{H}^{0} \right| \alpha\vec{R} \right\rangle$$

$$\begin{split} & \varepsilon_{\alpha\vec{R}} = \varepsilon_{\alpha\vec{R}}^{\quad 0} - e\Phi(\vec{R}) + \frac{1}{2}g_{0}\mu_{B}\vec{\sigma}\cdot\vec{B} & \text{Under}\\ & t_{\alpha\vec{R},\alpha'\vec{R}'} = t_{\alpha\vec{R},\alpha'\vec{R}'}^{\quad 0} \exp\left\{-\frac{ie}{2\hbar}\vec{B}\cdot\left(\vec{R}\times\vec{R}'\right)\right\} & \text{Symmetric gauge} \end{split}$$

P.Vogl, C. Strahberger phys. stat. sol. (b) 234, No. 1, 472. 477 (2002)





Distribution both of computational effort and memory resources.



% mexpensive+cluster of % desktop+computers

M.ZIELINSKI



Hard to code (320 cores!)



CRYSTAL LATTICE DISTRIBUTION



QNANO:

Parallel conjugate-gradient (VFF) Parallel Lanczos (TB) Distribution of TB matrix and TB matrix-vector product Parallel Coulomb ME Parallel fits (genetic alg.) Parallel LCAO

Including

- 1. Equal distribution (load balancing)
- 2. Proper (no deadlocks) and minimal (speed!) communication



M.ZIELINSKI

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STRAIN

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TIGHT BINDING METHOD EFFECT OF STRAIN

SURFACE PASSIVATION

EXTERNAL FIELDS

MANY-BODY EFFECTS

MULTI-EXCITON COMPLEXES CHARGED EXCITONS

EXAMPLES:

InAs/GaAs SELF-ASSEMBLED DOTS

(Zielinski et al PRB2010)

CdSe NANOCRYSTAL

REFERENCES





SAME QDOT, DIFFERENT STRAIN AND VALENCE BAND OFFSETS : DIFFERENT MODELS



SAME QDOT DIFFERENT METHODOLOGY



EBOM-effective bond orbital model, W.Sheng et al, PRB2005.

TB – Zielinski et al, PRB2010

EMP1(2) ó empirical pseudopotential results , He et al Phys. Rev. B 73, 115324 (2006)





IMS In-flush growth technique ó Wasilewski et al J.Cryst.Growth,1999



QD diameter 12.5 nm, height 2.0 nm VFF domain: 32.2 mln atoms TB domain: 0.298 mln atoms







QD diameter 16.0 nm, height from 5.0 nm to 2.0 nm VFF domain: 32.2 mln atoms TB domain: 0.44 mln atoms

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SINGLE PARTICLE SPECTRUM $E(m, n) = \Omega_{-}(m+1/2) + \Omega_{+}(n+1/2)$

FOCK-DARWIN SPECTRUM OF In FLUSH QDOTS



Raymond et al, PRL 2004, IMS/Grenoble/Dortmund



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INTERACTING ELECTRONS AND HOLES





INTERACTING ELECTRONS AND HOLES

HAMILTONIAN OF QUASI-ELECTRONS QUASI-HOLES:

$$\begin{split} \ddot{H} &= \sum_{i} E_{i}^{(e)} c_{i}^{+} c_{i}^{+} + \frac{1}{2} \sum_{ijkl} \langle i, j | V_{ee} | k, l \rangle c_{i}^{+} c_{j}^{+} c_{k} c_{l}^{\text{ELECTRON SP STATES,}} \\ &+ \sum_{j} E_{j}^{(h)} h_{j}^{+} h_{j}^{+} + \frac{1}{2} \sum_{ijkl} \langle i, j | V_{hh} | k, l \rangle h_{i}^{+} h_{j}^{+} h_{k} h_{l}^{\text{HOLE SP STATES,}} \\ &+ \sum_{j} E_{j}^{(h)} h_{j}^{+} h_{j}^{-} + \frac{1}{2} \sum_{ijkl} \langle i, j | V_{hh} | k, l \rangle h_{i}^{+} h_{j}^{+} h_{k} h_{l}^{\text{HOLE SP STATES,}} \\ &- \sum_{ijkl\sigma\sigma'} \langle \langle i_{e}, j_{h} | V_{eh,dir} | k_{h}, l_{e} \rangle - \langle \langle i_{e}, j_{h} | V_{eh,xchg} | l_{h}, k_{e} \rangle \rangle c_{i}^{+} h_{j}^{+} h_{k} c_{l} \end{split}$$

ELECTRON-HOLE DIRECT AND EXCHANGE INTERACTION

2.Build $\langle s | H | p \rangle$ 3.Diagonalize H, obtain eigenstates and eigenvalues

1. Build configurations $|p\rangle = c_i^+ c_j^+ h_k^+ h_l^+ ... |0\rangle$

EMISSION SPECTRA – FERMI'S GOLDEN RULE

$$I(\omega) = \sum_{f} \left| \left\langle f, N - 1 \middle| P^{-} \middle| i, N \right\rangle \right|^{2} \delta\left(E_{i} - E_{f} - \hbar \omega \right) \qquad P^{-} = \sum_{lm} \left\langle l_{e} \middle| \vec{\varepsilon} \cdot \vec{r} \middle| m_{h} \right\rangle c_{l} h_{m}$$

OPTICAL SPECTRA

$$\left\langle l_{e}\left|x\right|m_{h}\right\rangle = \sum_{R,\alpha} Ra_{R,\alpha}^{(e)*}a_{R,\alpha}^{(h)} + \sum_{R,\alpha} \sum_{\beta\neq\alpha} a_{R,\alpha}^{(e)*}a_{R,\beta}^{(h)}\left\langle \alpha\left|x\right|\beta\right\rangle + \sum_{R,\alpha} \sum_{R',\beta} a_{R,\alpha}^{(e)*}a_{R',\beta}^{(h)}\left\langle R,\alpha\left|x\right|R',\beta\right\rangle$$

ONSITE ONSITE NEAREST DIAGONAL OFFDIAGONAL NEIGHBORS



IN PRNCIPLE, WE SHOULD SOLVE FOR SCREENED COULOMB INTERACTIONS

$$\ddot{V}(r_1, r_2, \omega) = \ddot{V}^0(r_1, r_2) + \ddot{V}^0(r_1, r') \stackrel{\text{``I}}{\Pi}(r', r'', \omega) \stackrel{\text{''}}{V}(r'', r_2, \omega)$$
POLARIZATION OPERATOR

INTERACTIONS

BUT WE TAKE STATICALLY SCREENED INTERACTIONS

$$\ddot{V}(r_1, r_2) = \frac{e^2}{\varepsilon(r_1, r_2)|r_1 - r_2|}$$



 $\left|\Psi_{TB}^{e,h}\right\rangle = \sum_{atoms\vec{R}} \sum_{orbitals\,\alpha} c_{\vec{R}\alpha}^{e,h} \left|\vec{R}\,\alpha\right\rangle$

 $\left\langle \Psi_{i}\Psi_{j} \left| V \right| \Psi_{k}\Psi_{l} \right\rangle = \sum_{R_{1}\alpha_{1}} \sum_{R_{2}\alpha_{2}} \sum_{R_{3}\alpha_{3}} \sum_{R_{4}\alpha_{4}} a_{R_{1}\alpha_{1}}^{(i)*} a_{R_{2}\alpha_{2}}^{(j)*} a_{R_{3}\alpha_{3}}^{(k)} a_{R_{4}\alpha_{4}}^{(l)} \left\langle R_{1}\alpha_{1}, R_{2}\alpha_{2} \right| \frac{e^{2}}{\varepsilon |\vec{r_{1}} - \vec{r_{2}}|} \left| R_{3}\alpha_{3}, R_{4}\alpha_{4} \right\rangle$

COMPUTATIONAL EFFORT NONTRIVIAL!



$$\left\langle \Psi_{i}\Psi_{j}\left|V\right|\Psi_{k}\Psi_{l}\right\rangle = \sum_{R_{i}\alpha_{1}}\sum_{R_{2}\alpha_{2}}\sum_{R_{3}\alpha_{3}}\sum_{R_{4}\alpha_{4}}a_{R_{1}\alpha_{1}}^{(i)*}a_{R_{2}\alpha_{2}}^{(k)}a_{R_{3}\alpha_{3}}^{(l)}a_{R_{4}\alpha_{4}}^{(l)}\left\langle R_{1}\alpha_{1}, R_{2}\alpha_{2}\right|\frac{e^{2}}{\varepsilon|\vec{r_{1}}-\vec{r_{2}}|}|R_{3}\alpha_{3}, R_{4}\alpha_{4}\right\rangle$$

$$\left\langle \Psi_{i}\Psi_{j} \left| V \right| \Psi_{k}\Psi_{l} \right\rangle = \sum_{R_{1}} \sum_{\alpha_{1}\alpha_{2}\alpha_{3}\alpha_{4}} a_{R_{1}\alpha_{1}}^{(i)*} a_{R_{1}\alpha_{2}}^{(j)*} a_{R_{1}\alpha_{3}}^{(k)} a_{R_{1}\alpha_{4}}^{(l)} \left\langle R_{1}\alpha_{1}, R_{1}\alpha_{1} \right| \frac{e^{2}}{|\vec{r} - \vec{r}_{2}|} \left| R_{1}\alpha_{3}, R_{1}\alpha_{4} \right\rangle + \text{ON-SITE}$$

$$\sum_{R_{1},R_{NN}} \sum_{\alpha_{1}\alpha_{2}\alpha_{3}\alpha_{4}} a_{R_{1}\alpha_{1}}^{(i)*} a_{R_{1}\alpha_{1}}^{(j)*} a_{R_{1}\alpha_{4}}^{(l)} \left\langle R_{1}\alpha_{1}, R_{NN}\alpha_{2} \right| \frac{e^{2}}{|\vec{r}_{1} - \vec{r}_{2}|} \left| R_{NN}\alpha_{3}, R_{1}\alpha_{4} \right\rangle + \text{NN}$$

$$+ \sum_{R_{1}} \sum_{R_{2} \neq R_{1}} \left[\sum_{\alpha_{1}} a_{R_{1}\alpha_{1}}^{(i)*} a_{R_{1}\alpha_{1}}^{(l)} \right] \left[\sum_{\alpha_{2}} a_{R_{2}\alpha_{2}}^{(j)*} a_{R_{2}\alpha_{2}}^{(k)} \right] \frac{e^{2}}{|\vec{r}_{1} - \vec{R}_{2}|} \text{LONG-RANGE}$$



ELECTRON HOLE EXCHANGE

MONOPOLE-MONOPOLE, MONOPOLE-DIPOLE, DIPOLE-DIPOLE

$$\left\langle v'c' \middle| V^{X} \middle| vc \right\rangle = \sum_{R_{1} \neq R_{2}} \sum_{a_{1}b_{1}} \sum_{a_{2}b_{2}} (a_{R_{1}b_{1}}^{(v')}) (a_{R_{1}a_{1}}^{(c')})^{*} (a_{R_{2}b_{2}}^{(v)})^{*} (a_{R_{2}a_{2}}^{(c)}) \left\langle R_{1}a_{1}, R_{2}b_{2} \middle| \frac{1}{\left| \vec{r}_{1} - \vec{r}_{2} \right|} \middle| R_{2}a_{2}, R_{1}b_{1} \right\rangle = \iint dr_{1}dr_{2}\varphi_{b1}^{*} (r_{1} - R_{1})\varphi_{a1} (r_{1} - R_{1}) \frac{1}{\left| r_{1} - r_{2} \right|} \varphi_{b2} (r_{2} - R_{2})\varphi_{a2}^{*} (r_{2} - R_{2})$$
El-hole complex
El-hole complex
Output to P1

On atom at R1 On atom at R2



COULOMB MATRIX ELEMENTS ELECTRON HOLE EXCHANGE MONOPOLE-MONOPOLE, MONOPOLE-DIPOLE, DIPOLE-DIPOLE $\left\langle v'c' \middle| V^{X} \middle| vc \right\rangle = \sum_{R_{1} \neq R_{2}} \sum_{a_{1}b_{1}} \sum_{a_{2}b_{2}} \left(a_{R_{1}b_{1}}^{(v')}\right) \left(a_{R_{1}a_{1}}^{(c')}\right)^{*} \left(a_{R_{2}b_{2}}^{(v)}\right)^{*} \left(a_{R_{2}a_{2}}^{(c)}\right) \left\langle R_{1}a_{1}, R_{2}b_{2} \middle| \frac{1}{\left|\vec{r_{1}} - \vec{r_{2}}\right|} \middle| R_{2}a_{2}, R_{1}b_{1} \right\rangle$ $\left\langle R_{1}a_{1}, R_{2}b_{2} \left| \frac{1}{\left| \vec{r_{1}} - \vec{r_{2}} \right|} \right| R_{2}a_{2}, R_{1}b_{1} \right\rangle = \iint dr_{1}dr_{2}\varphi_{b1}^{*}(r_{1} - R_{1})\varphi_{a1}(r_{1} - R_{1}) \frac{1}{\left| r_{1} - r_{2} \right|} \varphi_{b2}(r_{2} - R_{2})\varphi_{a2}^{*}(r_{2} - R_{2})$ $\frac{1}{\left|\vec{r_1} - \vec{r_2}\right|} = \frac{1}{\left|\vec{R_1} - \vec{R_2}\right|} - \left[\left(\vec{r_1} - \vec{R_1}\right) - \left(\vec{r_2} - \vec{R_2}\right)\right] \frac{\left(\vec{R_1} - \vec{R_2}\right)}{\left|\vec{R_1} - \vec{R_2}\right|^3}$

 $+\sum_{i,j}(\vec{r_1}-\vec{R_1})_i(\vec{r_1}-\vec{R_1})_j(\frac{\left|\vec{R_1}-\vec{R_2}\right|^2\delta_{ij}-3(\vec{R_1}-\vec{R_2})_i(\vec{R_1}-\vec{R_2})_j}{\left|\vec{R_1}-\vec{R_2}\right|^5})+$



ELECTRON HOLE EXCHANGE MONOPOLE-MONOPOLE,í

$$\iint dr_1 dr_2 \,\varphi_{b1}^*(r_1 - R_1) \varphi_{a1}(r_1 - R_1) \frac{1}{|r_1 - r_2|} \varphi_{b2}(r_2 - R_2) \varphi_{a2}^*(r_2 - R_2)$$

COULOMB MATRIX ELEMENTS

Monopole-monopole contribution

$$\frac{1}{\left|\vec{r_1} - \vec{r_2}\right|} = \frac{1}{\left|\vec{R_1} - \vec{R_2}\right|} +$$

 $\frac{\delta_{a1b1}\delta_{a2b2}}{\left|\vec{R}_1-\vec{R}_2\right|} +$

This term is related to õsö state content in the hole state and õpö state content in the electron state

Other terms: dipole-monopole and dipole-dipole





INTERACTING ELECTRONS AND HOLES

EXAMPLE



MULTI-EXCITON COMPLEXES MULTI EXCITON COMPLEXES IN A BENCHMARK QD For pseudopotential calculations see L. He and A. Zunger Phys. Rev. B 73, 115324 (2006)

QD diameter 25nm, height 3.5 nm VFF domain: 50.3 mln atoms TB domain : 0.561 mln atoms







EXCITONIC ABSORPTION SPECTRUM







A $2E_{X}$ E_{XX} E_{X} E_{X} E_{X}

Benson et al

Real life






HIDDEN SYMMETRIES IN MULTI-EXCITON COMPLEXES





A.Wojs,PH, Solid State Comm. 100, 487 (1996).,P. H, Phys. Rev. B60, 5597 (1999); M.Bayer,I

HIDDEN SYMMETRIES ON DEGENERATE SHELLS

Interband polarisation operator
$$P^+ = \sum_i c_i^+ h_i^+$$

On shell $[H, P^+] = E_X P^+$
 $|N \rangle = (P^+)^N |0 \rangle \cdot E(N) = N E x$

SCALLE

multiplicative states=exciton condensate



$E (Q^{+} | 0 >) = E ((P^{+})^{2} | 0 >)$



EXCITONIC ARTIFICIAL ATOMS



A.Wojs,PH, Solid State Comm. 100, 487 (1996).,P. H, Phys. Rev. B60, 5597 (1999); M.Bayer,P. H, Nature 2000.

MULTIEXCITON EMISSION SPECTRA



QNANO:COMPUTATIONAL PLATFORM FOR ELECTRONIC PROPERTIES OF QUANTUM NANOSTRUCTURES

SUMMARY:

INTRODUCTION ATOMIC STRUCTURE DEFINITION STRAIN ELECTRONIC STRUCTURE CALCULATION TIGHT BINDING METHOD EFFECT OF STRAIN SURFACE PASSIVATION EXTERNAL FIELDS MANY-BODY EFFECTS MULTI-EXCITON COMPLEXES CHARGED EXCITONS EXAMPLES:

> InAs/GaAs SELF-ASSEMBLED DOTS CdSe NANOCRYSTAL

REFERENCES



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REFERENCES





CADMIUM SELENIDE

WURTZITE CRYSTAL STRUCTURE



INTERACTING ELECTRONS AND HOLES

ELECTRON-HOLE HAMILTONIAN

$$+\sum_{ijkl\sigma\sigma'} \langle i\sigma, j\sigma' | V | k\sigma', l\sigma \rangle c_{i\sigma}^{+} c_{j\sigma'} c_{k\sigma'} h_{l\sigma}$$

E-H NONCONSERVING TERMS

+
$$\sum_{ijkl\sigma\sigma'} \langle i\sigma, j\sigma' | V | k\sigma', l\sigma \rangle h_{i\sigma}^{+} h_{j\sigma'} c_{k\sigma'} h_{l\sigma} + \dots$$



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CdSe NANOCRYSTAL, GRAPHENE

REFERENCES:



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I.OZFIDAN, P.POTASZ. D. GUCLU - GRAPHENE

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