

### §3. The Cayley-Dickson process

The Cayley-Dickson process, as its name suggests, is due to A. A. Albert. It is a doubling process for building new algebras out of old ones. Given a scalar  $\mu \in \phi$  and a unital algebra  $B$  with scalar involution  $b \rightarrow \bar{b}$  (so  $b\bar{b} = \bar{b}b = n(b)$  and  $t(b) = b + \bar{b}$  lie in  $\phi 1$ ), we will build a new algebra

$$(3.1) \quad \mathbb{C}(B, \mu) = B \oplus B\ell$$

from two copies of  $B$  as linear space, with new involution

$$(3.2) \quad (b+c\ell)^* = \bar{b} - c\ell$$

and new multiplication given by the Cayley-Dickson formula

$$(3.3) \quad (b_1+c_1\ell)(b_2+c_2\ell) = (b_1b_2+\bar{c}_2c_1) + (c_2b_1+c_1\bar{b}_2)\ell.$$

This eminently forgettable formula can be broken down into bite-sized pieces. Besides the fact that  $B$  is imbedded as a subalgebra with its usual multiplication we have

$$(3.4) \quad b\ell = \ell\bar{b}$$

$$(3.5) \quad b(c\ell) = (cb)\ell$$

$$(3.6) \quad (c\ell)b = (c\bar{b})\ell$$

$$(3.7) \quad (b\ell)(c\ell) = \mu\bar{c}b.$$

Some helpful mnemonic devices: in (3.5) notice that to multiply  $b$  against  $c\ell$  you slip the  $b$  in behind the  $c$ , and also in (3.6) the  $b$  gets put between

the  $c$  and the  $\ell$ , but in moving past the  $\ell$  it gets conjugated as in (3.4). (WARNING: many authors use  $B \oplus \ell B$  instead, which turns all these formulas around.)

The multiplication formula is forced upon us: we must define multiplication by (3.3) if we want to have  $\ell^2 = \mu 1$  and  $b\ell = \ell \bar{b}$ .

3.8 (Necessity Proposition) If  $B$  is a unital subalgebra with scalar involution of a unital alternative algebra  $A$ , and  $\ell$  an element of  $A$  satisfying  $\ell^2 = \mu 1$  and  $b\ell = \ell \bar{b}$  for all  $b \in B$ , then  $B + B\ell$  is a subalgebra of  $A$  whose multiplication is given by the Cayley-Dickson formula.

Proof. Since (3.4) is one of our assumptions, we have  $(b\ell)(c\ell) = (\ell \bar{b})(c\ell) = \ell(\bar{b}c)\ell$  (Middle Moufang)  $= \ell\{\ell(\bar{b}c)\} = \ell^2(\bar{b}c) = \mu \bar{b}c$ , so (3.7) is forced upon us. From alternativity  $(cb)\ell + (c\ell)b = c(b\ell + \ell b)$  we see  $(c\ell)b = c\{(b + \bar{b})\ell\} - (cb)\ell$  (by (3.4))  $= \{c(b + \bar{b}) - cb\}\ell$  (since  $b + \bar{b} \in \phi 1$ )  $= (c\bar{b})\ell$ , so (3.6) too is forced. If  $A$  had an involution with  $\bar{\ell} = -\ell$  this would imply (3.5), but since we are not assuming an involution we argue dually  $b(c\ell) = b(\ell \bar{c}) = -\ell(b\bar{c}) + (b\ell + \ell b)\bar{c} = -\ell(b\bar{c}) + \{\ell(b + \bar{b})\}\bar{c} = \ell\{-b\bar{c} + (b + \bar{b})\bar{c}\} = \ell(\bar{b}\bar{c}) = \ell(\overline{cb}) = (bc)\ell$ .

If  $\mu$  is cancellable ( $\mu x = 0 \Rightarrow x = 0$ ) we could also derive (3.5) - (3.6) from the Moufang formulas: for example,  $\mu\{b(c\ell)\} = \ell\{\ell\{b(\ell \bar{c})\}\} = \ell\{(\ell b\ell)\bar{c}\} = \ell\{(\ell^2 \bar{b})\bar{c}\} = \mu \ell(\bar{b}\bar{c}) = \mu(cb)\ell$  and  $\mu\{(c\ell)b\} = \{(c\ell)b\}\ell = (c\{\ell b\ell\})\ell = \mu(c\bar{b})\ell$ .

Thus the various pieces (3.4) - (3.7) of the Cayley-Dickson formula are forced upon us. In particular,  $B + B\ell$  is a subalgebra.  $\square$

3.9 Corollary. If  $B$  is a unital subalgebra with scalar involution of a unital alternative algebra  $A$  and  $\ell \in A$  an element satisfying  $\ell^2 = \mu 1$  and  $b\ell = \ell \bar{b}$ , then the map  $b \oplus c\ell \rightarrow b + c\ell$  is a homomorphism  $\mathbb{C}(B, \mu)$  onto the subalgebra  $B + B\ell$ . When  $A$  has an involution extending that on  $B$  with  $\ell^* = -\ell$ , this homomorphism is a  $*$ -homomorphism. When  $B \cap B\ell = 0$  and  $\mu$  is cancellable, the map is an isomorphism.

Proof. Since both  $\mathbb{C}(B, \mu)$  and  $B + B\ell$  have multiplication given by the Cayley-Dickson formula, the map is a homomorphism. If  $B \cap B\ell = 0$  then  $b + c\ell = 0 \Rightarrow b = c\ell = 0$ , and if  $\mu$  is cancellable  $c\ell = 0 \Rightarrow \mu c = (c\ell)\ell = 0 \Rightarrow c = 0$ , so in these cases the map is injective. If an involution  $*$  on  $A$  satisfies  $\ell^* = -\ell$  and  $b^* = \bar{b}$ , then  $b\ell = \ell \bar{b} \Rightarrow b\ell = -\ell^*b^* = -(b\ell)^* \Rightarrow (b\ell)^* = -b\ell$ , so the involution on  $A$  satisfies (3.2). In this case the map preserves the involution.  $\square$

We say an algebra  $\mathbb{C}(B, \mu)$  (alternative or not) is obtained from  $B$  by the Cayley-Dickson process. Note that this process doubles dimension.

Returning to  $\mathbb{C}(B, \mu)$  in general, let us verify that  $*$  is again a scalar involution. Clearly it is linear of period 2. It is an anti-homomorphism since for  $x_1 = b_1 + c_1\ell$  we have  $\overline{x_1 x_2} = \overline{(b_1 b_2 + \mu c_2 c_1) + (c_2 b_1 + c_1 \bar{b}_2)\ell} = (\bar{b}_2 \bar{b}_1 + \mu \bar{c}_1 \bar{c}_2) - (c_1 \bar{b}_2 + c_2 \bar{b}_1)\ell = (\bar{b}_2 - c_2\ell)(\bar{b}_1 - c_1\ell) = \overline{x_2 x_1}$  by (3.3). Furthermore, the involution is scalar since

$$\begin{aligned} \text{tr}(x) &= \text{tr}(b) \\ (3.10) \quad n(x) &= n(b) - \mu n(c) \end{aligned} \quad (x = b + c\ell)$$

The trace is easy, while for the norm we have  $x\bar{x} = (b + c\ell)(\bar{b} - c\ell) = b\bar{b} - \mu c\bar{c} = n(b) - \mu n(c)$ .

Notice that since it possesses a scalar involution,  $\mathbb{C}(B, \mu)$  is of degree 2.

From the expression for the norm on  $\mathbb{C}(B, \mu)$ , we see that it inherits nondegeneracy from  $B$ :

3.11 (Nondegeneracy Criterion) If  $\mu$  is cancellable then  $n(x, y)$  is nondegenerate as a bilinear form on  $\mathbb{C}(B, \mu)$  if it is nondegenerate on  $B$ . If in addition  $n(b) \neq \mu n(c)$  whenever  $n(b), n(c) \neq 0$  but  $n(b, B) = n(c, B) = 0$ , then  $n$  is nondegenerate as a quadratic form on  $\mathbb{C}(B, \mu)$  if it is nondegenerate on  $B$ .

Proof. If  $z = b + c\lambda$  has  $n(z, x) = 0$  for all  $x = a + d\lambda$  then from linearized (3.10)  $n(b, a) - \mu n(c, d) = 0$  for all  $a, d$ ; setting successively  $d = 0$ ,  $a = 0$  we see  $n(b, B) = \mu n(c, B) = 0$ . If  $n(x, y)$  is nondegenerate on  $B$  and  $\mu$  is cancellable, this yields  $b = c = 0$  and  $z = 0$ . If we only assume the quadratic form is nondegenerate on  $B$ , then when  $n(z) = n(z, x) = 0$  for all  $x$  we have  $n(b) = \mu n(c)$  yet  $n(b, B) = n(c, B) = 0$ , hence  $n(b), n(c) \neq 0$  by nondegeneracy on  $B$ , contrary to our hypothesis  $n(b) \neq \mu n(c)$ .  $\square$

3.12 Remark. If  $\phi$  is a field, the condition  $n(b) \neq \mu n(c)$  is equivalent to the condition that  $\mu \neq n(b)n(c)^{-1} = n(bc^{-1})$  not be a norm,  $\mu \notin n(B)$ .  $\square$

To see what sort of algebras the Cayley-Dickson process leads to, we prove

3.13 (Criterion) If  $B$  is a unital algebra with scalar involution then  
(1)  $\mathbb{C}(B, \mu)$  is commutative iff  $B$  is commutative with trivial

involution; (2)  $\mathbb{C}(B, \mu)$  is associative iff  $B$  is commutative and associative; (3)  $\mathbb{C}(B, \mu)$  is alternative iff  $B$  is associative.

Proof. From the Cayley-Dickson formula the left and right multiplications by an element  $x = b + c\bar{c}$  in  $\mathbb{C}(B, \mu)$  have matrices

$$L_x = \begin{pmatrix} L_b & \mu R_c^* \\ L_c^* & R_b \end{pmatrix} \quad R_x = \begin{pmatrix} R_b & \mu L_{\bar{c}} \\ L_c & R_{\bar{b}} \end{pmatrix}$$

relative to the decomposition  $\mathbb{C}(B, \mu) = B \oplus B\bar{c}$ :  $L_x b_2 = (bb_2) \oplus (c\bar{b}_2)\bar{c}$   
 $= \{L_b(b_2)\} \oplus \{L_c^*(b_2)\}\bar{c}$ ,  $L_x(c_2\bar{c}) = \mu\bar{c}_2c \oplus (c_2b)\bar{c} = \{\mu R_c^*(c_2)\} \oplus$   
 $\{R_b(c_2)\}\bar{c}$ ,  $R_x b_1 = (b_1b) \oplus (cb_1)\bar{c} = \{R_b(b_1)\} \oplus \{L_c(b_1)\}\bar{c}$ ,  $R_x(c_1\bar{c}) =$   
 $\mu\bar{c}c_1 \oplus (c_1\bar{b})\bar{c} = \{\mu L_{\bar{c}}(c_1)\} \oplus \{R_{\bar{b}}(c_1)\}\bar{c}$ . We will deal mainly with the  
 $R_x$ , since they have a slightly simpler form.

Commutativity of  $\mathbb{C}$  means  $L_x = R_x$  for all  $x$ , and this is equivalent to  $L_b = R_b$  and (setting  $c = 1$ ,  $L_c = 1$ )  $\mu = I$ , i.e. that  $B$  is commutative with identity involution.

Associativity of  $\mathbb{C}$  means  $R_{x_2} R_{x_1} = R_{x_1 x_2}$ , i.e.

$$\begin{pmatrix} R_{b_2} R_{b_1} + L_{\bar{c}_2} L_{c_1} & \mu \{R_{b_2} L_{\bar{c}_1} + L_{\bar{c}_2} R_{b_1}\} \\ L_{c_2} R_{b_1} + R_{b_2} L_{c_1} & \mu L_{c_2} L_{\bar{c}_1} + R_{b_2} R_{b_1} \end{pmatrix} = \begin{pmatrix} R_{b_1 b_2} + R_{c_2 c_1} & \mu \{L_{b_1} \bar{c}_2 + L_{b_2} \bar{c}_1\} \\ L_{c_2 b_1} + L_{c_1 b_2} & \mu R_{b_2} \bar{c}_1 + \mu R_{c_1} c_2 \end{pmatrix}.$$

Setting  $c_1 = c_2 = 0$  shows  $R_{b_2} R_{b_1} = R_{b_1 b_2}$ , i.e.  $B$  is associative; setting  $b_2 = c_1 = 0$  and  $c_2 = 1$  shows  $R_{b_1} = L_{b_1}$  (lower left entry), i.e.  $B$  is commutative. Conversely, if  $B$  is commutative associative then  $R_b R_c = R_{cb}$ ,

$L_b L_c = L_{bc}$ ,  $L_b = R_b$  show  $R_{x_2} R_{x_1} = R_{x_1 x_2}$ .

Alternativity of  $\mathbb{C}$  means  $R_x^2 = R_x^2$  (the involution then gives left alternativity), or equivalently  $R_x R_{\bar{x}} = R_{x\bar{x}}$  (since  $x + \bar{x} = t(x) \in \phi 1$ ),

$$\begin{pmatrix} R_b R_{\bar{b}} - \mu L_{\bar{c}} L_c & \mu \{-R_b L_{\bar{c}} + L_{\bar{c}} R_b\} \\ L_c R_{\bar{b}} - R_b L_c & -\mu L_{\bar{c}} L_{\bar{c}} + R_b R_b \end{pmatrix} = (n(b) - \mu n(c)) \begin{pmatrix} I & 0 \\ 0 & I \end{pmatrix}.$$

From  $[L_c, R_{\bar{b}}] = 0$  we see  $B$  is associative. Conversely, if  $B$  is associative the off-diagonals vanish and  $R_b R_{\bar{b}} - \mu L_{\bar{c}} L_c = R_{b\bar{b}} - \mu L_{\bar{c}c} = (n(b) - \mu n(c))I = R_b R_{\bar{b}} - \mu L_c L_{\bar{c}}$ , proving  $R_x R_{\bar{x}} = R_{x\bar{x}}$ .  $\square$

Thus you must start with an associative algebra in order to obtain an alternative one by the Cayley-Dickson construction; once you have built an alternative but not associative algebra you can go no further, anything further will not be alternative.

We are now in a position to build composition algebras. We begin with  $B_1$  being just the ring  $\phi 1$  of scalars with identity involution. Next we form a two dimensional  $B_2 = \mathbb{C}(B_1, \mu_1) = \phi 1 + \phi i$ , which will be commutative associative since  $B_1$  is commutative associative with identity involution; if  $\phi$  has characteristic  $\neq 2$ ,  $2 \neq 0$ , the involution  $\alpha + \beta i \rightarrow \alpha - \beta i$  is not the identity.

In characteristic 2 the usual process applied to any algebra with identity involution will still have identity involution. To break out of this cycle we must modify the process slightly. We form

$$B'_2 = \phi 1 + \phi u \quad u^2 - u + \mu_1 1 = 0$$

so that  $B_2$  is a commutative associative algebra with nontrivial involution determined by

$$u + \bar{u} = 1.$$

Indeed, the only non-trivial product of basis elements 1,  $u$  is  $u^2$ , and  $\bar{u}^2 = (1-u)^2 = 1 - 2u + u^2 = 1 - u - \mu_1 1 = \bar{u} - \mu_1 1 = \overline{u - \mu_1 1} = \overline{u^2}$ . Note  $t(u) = 1$ ,  $n(u) = -\mu_1$ . We will still call this algebra  $\mathbb{C}(B_1, \mu_2)$  and pretend it is obtained by the Cayley-Dickson process. Remember: the second stage of the Cayley-Dickson process is defined differently in characteristic 2.

Once we have arrived at a second-stage algebra  $B_2$  of dimension 2 with nontrivial scalar involution, we can form a 4-dimensional  $B_3 = \mathbb{C}(B_2, \mu_2) = B_2 \oplus B_2 j$  which will be associative but not commutative. These are precisely the quaternion algebras over  $\phi$ .

From  $B_3$  we construct an 8-dimensional algebra  $B_4 = \mathbb{C}(B_3, \mu_3)$  which will be alternative but (since  $B_3$  is not commutative) not associative. Such an 8-dimensional algebra is called a Cayley algebra (alias Cayley-Dickson, Cayley-Graves, Albert-Dickson, octaves, you name it; in analogy with quaternions, they are often called octonions). Thus Cayley algebras are those obtained from quaternion algebras by the Cayley-Dickson process.

At this point the construction stops, for further algebras will not be alternative.

These algebras

$$B_1 = \phi 1 \quad B_2 = \mathbb{C}(B_1, \mu_1) \quad B_3 = \mathbb{C}(B_2, \mu_2) \quad B_4 = \mathbb{C}(B_3, \mu_3)$$

of dimensions 1, 2, 4, 8 are the basic algebras obtained by the Cayley-

Dickson process, and so we will call them the Cayley-Dickson process algebras. If we had let the general construction run on in Characteristic 2 (rather than modifying it) we would simply have gotten larger and larger commutative, associative algebras with identity involution such that  $x^2 \in \phi$  for all  $x$ . If  $\phi$  is a field, this means we have a purely inseparable extension of exponent 2.

As alternative algebras with scalar involution, the Cayley-Dickson process algebras are of degree 2 and their norm forms permit composition (see (2.4)). To be composition algebras according to our definition, the norms must be nondegenerate. In the case of  $B_1$ ,  $n(\alpha, \beta) = 2\alpha\beta$  is nondegenerate iff  $\phi$  has no 2-torsion, and  $n(\alpha) = \alpha^2$  is nondegenerate iff  $\phi$  has no elements with  $\alpha^2 = 2\alpha = 0$ . If  $\phi$  has no 2-torsion and we take  $\mu_1, \mu_2, \mu_3$  to be cancellable, then by the Nondegeneracy Criterion 3.11 the quadratic extension  $B_2$ , the quaternion algebra  $B_3$ , and the Cayley algebra  $B_4$  will all have nondegenerate forms  $n(x, y)$ , and so are composition algebras. If  $\phi$  has characteristic 2 the form  $n(x, y)$  on the modified  $C_1(\phi, \mu_1) = \phi 1 + \phi u$  is nondegenerate (being hyperbolic,  $n(\alpha 1 + \beta u, \alpha' 1 + \beta' u) = \alpha\beta' + \alpha'\beta$  if  $2 = 0$ ), so again  $B_2, B_3, B_4$  are composition algebras.

3.14 (Composition Proposition) If  $\phi$  has no 2-torsion or all 2-torsion, then any algebra of dimension 2, 4 or 8 obtained by the Cayley-Dickson process by means of cancellable parameters  $\mu_1, \mu_2, \mu_3$  is a composition algebra with nondegenerate norm bilinear form  $n(x, y)$ ; degeneracy of  $n(x, y)$  is possible only in dimension 1. If in addition  $\phi$  has no nilpotent elements, then the norm form  $n$  is nondege-



nerate in all cases and all the algebras are composition algebras.  $\square$

In the next section we will see that conversely, over a field, all composition algebras are obtained by the Cayley-Dickson process.

In case  $\phi$  has no 2-torsion, the algebras we have built look like the following.

#### Dimension 1: Base ring

$B_1 = \phi 1 = \phi$  is a commutative associative algebra with identity involution  $\bar{1} = 1$  and norm  $n(\alpha_0) = \alpha_0^2$ .

#### Dimension 2: Quadratic extension

$B_2 = \phi 1 \oplus \phi i$  is a commutative associative algebra with nontrivial involution  $\bar{i} = -i$ , multiplication  $i^2 = \mu_1 1$ , and norm  $n(\alpha_0 1 + \alpha_1 i) = \alpha_0^2 - \alpha_1^2 \mu_1$ .

#### Dimension 4: Quaternion algebra

$B_3 = \phi 1 \oplus \phi i \oplus \phi j \oplus \phi k$  is a non-commutative associative algebra with involution  $\bar{i} = -i$ ,  $\bar{j} = -j$ ,  $\bar{k} = -k$ , norm

$$n(\alpha_0 1 + \alpha_1 i + \alpha_2 j + \alpha_3 k) = \alpha_0^2 - \alpha_1^2 \mu_1 - \alpha_2^2 \mu_2 + \alpha_3^2 \mu_1 \mu_2,$$

and multiplication

$$i^2 = \mu_1 1 \quad j^2 = \mu_2 1 \quad k^2 = -\mu_1 \mu_2 1$$

$$ij = k = -ji \quad jk = -\mu_2 i = -kj \quad ki = -\mu_1 j = -ik.$$

#### Dimension 8: Cayley algebra

$B_4 = \phi 1 \oplus \phi i \oplus \phi j \oplus \phi k \oplus \phi l \oplus \phi il \oplus \phi jl \oplus \phi kl$  is an alternative but not associative algebra with involution  $\bar{i} = -i, \bar{j} = -j, \bar{k} = -k, \bar{l} = -l, \overline{(il)} =$

$$-i\ell, \overline{(j\ell)} = -j\ell, \overline{(k\ell)} = -k\ell, \text{ norm}$$

$$\begin{aligned} n(\alpha_0 1 + \alpha_1 i + \alpha_2 j + \alpha_3 k + \alpha_4 \ell + \alpha_5 i\ell + \alpha_6 j\ell + \alpha_7 k\ell) \\ = \alpha_0^2 - \alpha_1^2 \mu_1 - \alpha_2^2 \mu_2 + \alpha_3^2 \mu_1 \mu_2 - \alpha_4^2 \mu_3 + \alpha_5^2 \mu_1 \mu_3 + \alpha_6^2 \mu_2 \mu_3 - \alpha_7^2 \mu_1 \mu_2 \mu_3 \end{aligned}$$

and multiplication

$$i^2 = \mu_1 1 \quad j^2 = \mu_2 1 \quad k^2 = -\mu_1 \mu_2 1 \quad \ell^2 = \mu_3 1 \quad (i\ell)^2 = -\mu_1 \mu_3 1 \quad (j\ell)^2 = -\mu_2 \mu_3$$

$$(k\ell)^2 = \mu_1 \mu_2 \mu_3 1 \quad ij = k = -ji \quad jk = -\mu_2 i = -kj \quad ki = -\mu_1 j = -ik$$

$$\ell i = -i\ell \quad \ell j = -j\ell \quad \ell k = -k\ell$$

$$i(i\ell) = \mu_1 \ell = -(i\ell)i \quad j(j\ell) = \mu_2 \ell = -(j\ell)j \quad k(k\ell) = -\mu_1 \mu_2 \ell = -(k\ell)k$$

$$i(j\ell) = -k\ell = -(j\ell)i \quad j(k\ell) = \mu_2 i\ell = -(k\ell)j \quad k(i\ell) = \mu_1 j\ell = -(i\ell)k$$

$$j(i\ell) = k\ell = -(i\ell)j \quad k(j\ell) = -\mu_2 i\ell = -(j\ell)k \quad i(k\ell) = -\mu_1 j\ell = -(k\ell)i$$

$$(i\ell)(j\ell) = \mu_3 k = -(j\ell)(i\ell) \quad (j\ell)(k\ell) = -\mu_2 \mu_3 i = -(k\ell)(j\ell) \quad (k\ell)(i\ell) =$$

$$-\mu_1 \mu_3 j = -(i\ell)(k\ell)$$

$$(i\ell)\ell = \mu_3 i = -\ell(i\ell) \quad (j\ell)\ell = \mu_3 j = -\ell(j\ell) \quad (k\ell)\ell = \mu_3 k = -\ell(k\ell)$$

The results in characteristic 2 are (even?) less memorable, since in place of  $1, i$  we have  $1, u$  with  $u^2 = u + \mu_2 1$ . Multiplication in the Cayley algebra can be summarized in a multiplication table for the products  $xy$ :

x \ y	1	i	j	k	ℓ	iℓ	jℓ	kℓ
1	1	i	j	k	ℓ	iℓ	jℓ	kℓ
i	i	$\mu_1 1$	k	$\mu_1 j$	iℓ	$\mu_1 \ell$	-kℓ	$-\mu_1 j \ell$
j	j	-k	$\mu_2 1$	$-\mu_2 i$	jℓ	kℓ	$\mu_2 \ell$	$\mu_2 i \ell$
k	k	$-\mu_1 j$	$\mu_2 i$	$-\mu_1 \mu_2 1$	kℓ	$\mu_1 j \ell$	$-\mu_2 i \ell$	$-\mu_1 \mu_2 \ell$
ℓ	ℓ	-iℓ	-jℓ	-kℓ	$\mu_3 1$	$-\mu_3 i$	$-\mu_3 j$	$-\mu_3 k$
iℓ	iℓ	$-\mu_1 \ell$	-kℓ	$-\mu_1 j \ell$	$\mu_3 i$	$-\mu_1 \mu_3 1$	$\mu_3 k$	$\mu_1 \mu_3 j$
jℓ	jℓ	kℓ	$-\mu_2 \ell$	$\mu_2 i \ell$	$\mu_3 j$	$-\mu_3 k$	$-\mu_2 \mu_3 1$	$-\mu_2 \mu_3 i$
kℓ	kℓ	$\mu_1 j \ell$	$-\mu_2 i \ell$	$\mu_1 \mu_2 \ell$	$\mu_3 k$	$-\mu_1 \mu_3 j$	$\mu_2 \mu_3 i$	$\mu_1 \mu_2 \mu_3 1$

Staring at this table usually is NOT the best way to understand the structure of a Cayley algebra.

If  $\phi = \mathbb{R}$  is the field of real numbers and  $\mu_1 = \mu_2 = \mu_3 = -1$  then the quadratic extension  $B_2$  is simply the field  $\mathbb{C}$  of complex numbers with conjugation as involution,  $B_3$  is the division algebra  $\mathbb{Q}$  of (ordinary) quaternions with standard involution, and  $B_4$  is the algebra  $\mathcal{L}$  of Cayley numbers; it is an alternative division algebra. It is often helpful to think of the Cayley-Dickson process algebras as generalizations of

$$\mathbb{R}, \mathbb{C}, \mathbb{Q}, \mathcal{L}.$$

We have previously referred to the following result in connection with one-sided and quadratic ideals:

3.15 (One-Sided Ideal Proposition) A Cayley algebra (split or not) over a field  $\Phi$  contains no proper one-sided ideals.

Proof. Because of the symmetry resulting from the involution, it suffices to show there are no proper left ideals  $B$  in a Cayley algebra  $\mathbb{C} = \mathbb{C}(A, \mu)$  ( $A$  a quaternion algebra). If  $B$  is nonzero it contains an element  $c + b\lambda$  for  $b \neq 0$  (if  $c + 0\lambda \in B$  then also  $\lambda(c + 0\lambda) = \lambda c = \bar{c}\lambda \in B$  for  $\bar{c} \neq 0$ ); since it is a left ideal,  $B$  also contains  $z\{(xy)(c+b\lambda) - x[y(c+b\lambda)]\} = z\{(xy)c - x(yc)\} + \{[b(xy) - (by)x]z\}\lambda = \{b[x, y]z\}\lambda$  for all  $x, y, z \in A$ . Thus  $B$  contains  $\{b[A, A]A\}\lambda = \{bA\}\lambda$  ( $[A, A]A$  is an ideal in the simple quaternion algebra  $A$ ). So far we have left-multiplied  $B$  only by  $A$ . If we now left-multiply by  $A\lambda$  we see  $B$  contains  $A\lambda\{(bA)\lambda\} = \mu(\bar{A}\bar{b})A = A\bar{b}A$ . Again by simplicity,  $\bar{b} \neq 0$  implies  $B$  contains  $A$ , therefore  $\lambda \cdot A = \bar{A}\lambda = A\lambda$  as well, and  $B = A + A\lambda = \mathbb{C}$ . Thus as soon as  $B \neq 0$  we have  $B = \mathbb{C}$ .  $\square$

3.16 Corollary. A Cayley algebra over a field is a simple alternative algebra of dimension 8.  $\square$

By means of the Inverse Criterion 2.7 for composition algebras we can decide when the Cayley-Dickson process yields a division algebra.

3.17 (Division Algebra Construction) If  $B$  is a composition algebra over a field then  $\mathbb{C}(B, \mu)$  is a division algebra iff  $B$  is a division algebra and  $\mu \notin n(B)$  is not a norm.

Proof. The condition  $\mathbb{C}(B, \mu)$  be a division algebra is that its norm not represent zero,  $n(x) = n(b) - \mu n(c) \neq 0$  for  $x = b + c\lambda \neq 0$ . Clearly

this implies  $n(b) \neq 0$  for  $b \neq 0$  (take  $c = 0$ ), so  $B$  must be a division algebra, and  $n(b) \neq \mu$  (take  $c = 1$ ), so  $\mu$  is not a norm.

Conversely, if these two conditions are met then  $n(b) - \mu n(c) \neq 0$  if  $c = 0$  but  $b \neq 0$  (since  $n(b) \neq 0$ ) and  $n(b) - \mu n(c) \neq 0$  if  $c \neq 0$  (since  $\mu \neq n(b)n(c)^{-1} = n(bc^{-1})$ ), so  $n(x) \neq 0$  for  $x \neq 0$ .  $\square$

On several occasions we will need invertible elements of various forms.

3.18 Lemma. Any Cayley-Dickson process algebra of dimension  $\geq 2$  contains cancellable skew elements  $x - \bar{x}$ , and in dimension  $\geq 4$  contains cancellable commutators  $[x, y]$ . If  $\phi$  is a field all these elements are invertible. If  $\phi$  is a field with more than 2 elements, the algebra has a basis of invertible elements.

Proof. When  $\phi$  has no 2-torsion, the element  $i - \bar{i} = 2i$  in  $B_2$  is cancellable since  $\mu_1$  is:  $ix = 0 \Rightarrow \mu_1 x = i(ix) = 0 \Rightarrow x = 0$ . In characteristic 2, the element  $u - \bar{u} = u + \bar{u} = 1$  is invertible in  $B_2'$ .

In characteristic  $\neq 2$  the commutator  $[i, j] = ij - ji = 2ij = 2k$  in  $B_3$  is cancellable since  $n(k) = + \mu_1 \mu_2$  is. In characteristic 2  $[u, j] = uj - ju = (u - \bar{u})j = (u + \bar{u})j = j$  is cancellable.

If  $\phi$  is a field then  $x$  cancellable  $\Rightarrow n(x) \neq 0 \Rightarrow n(x)$  is invertible  $\Rightarrow x$  is invertible by the Inverse Criterion 2.7.

If  $B$  has a basis of invertible  $\{b_i\}$  then  $\mathbb{C}(B, \mu)$  has a basis of invertible elements  $\{b_i\}$  and  $\{b_i \ell\}$ , and we start off with  $B_1 = \phi \ell$  having invertible basis  $\{\ell\}$ . (In characteristic 2,  $B_2' = \phi \ell + \phi u$  for  $u^2 - u + v \ell = 0$  has invertible basis  $\{\ell, \lambda \ell + u\}$  as long as  $0 \neq n(\lambda \ell + u) = \lambda^2 + \lambda t(u) + n(u) = \lambda^2 + \lambda + v$ ; since  $\phi$  has at least 3 elements and  $\lambda^2 + \lambda + v = 0$  at

most 2 roots, there is at least one non-root  $\lambda$ .)  $\square$

Let us also observe that the symmetric elements of a Cayley-Dickson process algebra are just those in  $\phi 1$ , except when the characteristic is 2 and the dimension is 4 or 8. This is clear from (3.2) in characteristic  $\neq 2$ , and also in characteristic 2 for dimension 1. In characteristic 2, dimension 2 the only element  $x = \beta u + \alpha \bar{u}$  which is symmetric,  $x = \bar{x} = \beta u + \alpha \bar{u}$ , is  $x = \alpha u + \alpha \bar{u} = \alpha 1$ .

Clearly everything in a Cayley-Dickson process algebra of dimension 1, 2, 4 associates and in dimension 1, 2 everything commutes. The unclear cases are clarified by

3.19 (Commuting and Associating Criterion) An element of a quaternion algebra commutes with everything iff it is a scalar in  $\phi 1$ . An element of a Cayley algebra commutes with everything or associates with everything iff it is a scalar in  $\phi 1$ .

Proof. If  $z = b + cj$  is in the center of  $\mathbb{O} = \mathbb{C}(B_2, \mu_2)$  then  $0 = [a, z] = [a, b] + (ca - \bar{c}\bar{a})j$  for all  $a \in B_2$  implies  $c(a - \bar{a}) = 0$ ; since there are invertible elements in  $B_2$  of the form  $a - \bar{a}$  by the Lemma 3.18, this implies  $c = 0$  and  $z = b$ . Then  $0 = [z, j] = (b - \bar{b})j$  implies  $b = \bar{b}$  is a scalar in  $B_2$ ,  $b = \beta 1$ , and  $z = \beta 1$ .

If  $z = b + c\ell$  commutes with everything in  $\mathbb{C} = \mathbb{C}(B_3, \nu_3)$  then  $0 = [a, z] = [a, b] + \{c(a - \bar{a})\}\ell$  shows  $[a, b] = c(a - \bar{a}) = 0$  for all  $a \in B_3$ . By the previous result, if  $b$  lies in the center of  $B_3$  we have  $b = \beta 1$ , and again since invertible  $a - \bar{a}$  exist in  $B_3$  we have  $c = 0$ , so  $z = \beta 1$ .

If  $z$  associates with everything in  $\mathbb{C}$  then  $0 = [a, d, z] = [a, d, b] + \{c(ad) - (cd)a\}z = \{c[a, d]\}z$  implies  $c[a, d] = 0$  for all  $a, d$  in the associative algebra  $Q = B_3$ . Once more by the Lemma 3.18,  $B_3$  has invertible commutators  $[a, d]$ , so  $c$  must be zero. Then  $0 = [a, z, z] = [a, b, z] = (ab - ba)z$  for all  $a$  implies  $b = \beta 1$  lies in the center of  $Q$ , so once more  $z = \beta 1$ .  $\square$

3.20 Corollary. If an element of a Cayley-Dickson process algebra commutes with everything, it also associates with everything.  $\square$

3.21 Corollary. A Cayley-Dickson process algebra is either associative, or else every element which associates with everything also commutes with everything.  $\square$

These properties will turn out to be general properties of alternative algebras. Commutativity will force associativity (except in characteristic 3; see Section III.4), and for simple not-associative algebras nucleus and center will coincide (see III.1). This is important for the general structure theory (see Appendix II).

## Exercise

- 3.1 Find the Cayley-Dickson formula for multiplication in  $B + \ell B$  if  $\ell^2 = \mu 1$ ,  $\ell b = \bar{b}\ell$ .
- 3.2 Try to define a general "modified Cayley-Dickson process"  $\mathbb{C}(B, v) = Bu \oplus B\bar{u}$  with  $u + \bar{u} = 1$ ,  $u^2 - u + v1 = 0$ .
- 3.3 In the "modified" construction  $B'_2 = \phi 1 \oplus \phi u$  in arbitrary characteristic,  $u^2 - u + v1 = 0$ , show  $n(x, y)$  is nondegenerate iff  $1 + 4v$  is cancellable. Conclude in characteristic 2 it is nondegenerate no matter what  $v$  is chosen.
- 3.4 If  $\frac{1}{2} \in \phi$  show  $u = \frac{1}{2}(1 + i)$  in  $B_1 = \mathbb{C}(\phi, \mu_1) = \phi 1 \oplus \phi i$  has  $u + \bar{u} = 1$ ,  $u^2 - u + v_1 1 = 0$ . What is  $1 + 4v_1$ ?
- 3.5 If  $B$  is a simple algebra with involution which is associative but not commutative, and  $\mu \in \phi$  is invertible, show the algebra  $\mathbb{C}(B, \mu)$  defined as in (3.1) with the Cayley-Dickson formula (3.3) has no proper ideals, in particular is simple. (Notice  $B$  is not assumed unital nor the involution scalar, so  $\mathbb{C}(B, \mu)$  need not be alternative. Also the element  $\ell$  need not exist in  $\mathbb{C}$ , so  $b\ell$  cannot be interpreted as a product of  $b$  with  $\ell$ .)
- 3.6 If  $B$  is unital but the involution not necessarily scalar, show for cancellable  $\mu$  that  $\mathbb{C}(B, \mu)$  is commutative iff  $B$  is commutative and  $* = 1$  (examine commutators); show  $\mathbb{C}$  is associative iff  $B$  is commutative associative (examine associators); show  $\mathbb{C}$  is alternative iff (i)  $B$  is alternative, (ii) all  $n(b) = b\bar{b} = \bar{b}b$  commute with  $B$ , (iii) all  $b + t(b)$  associate with  $B$  (examine associators). Show from (iii) that if  $B$  has no 3-torsion then  $B$  is associative, so  $*$  is a central



involution. Thus to get an alternative  $\mathbb{C}$  the involution has to be at least central anyway; regarding  $B$  as an algebra over its center  $\Omega$ , this means  $*$  is scalar over  $\Omega$ .

- 3.7 Give an example of a 2-dimensional Cayley-Dickson process algebra over the field  $\mathbb{Z}_2$  which has no basis of invertible elements. Can you give a 4-dimensional example?