

Phase Transition in Dirac Ensembles

Masoud Khalkhali

University of Western Ontario

Joint works with S. Azarfar, N. Pagliaroli, H. Hessam

Metric from Dirac operators

- ▶ **Connes' distance formula:** Let (M, g) be a spin Riemannian manifold with Dirac operator D Then

$$d(p, q) = \text{Sup}\{|f(p) - f(q)|; \|[D, f]\| \leq 1\}.$$

Compare with the classical dual formula

$$d(p, q) = \text{Inf} \int_p^q (g_{\mu\nu} dx^\mu dx^\nu)^{1/2}$$

Metric from Dirac operators

- ▶ **Connes' distance formula:** Let (M, g) be a spin Riemannian manifold with Dirac operator D . Then

$$d(p, q) = \text{Sup}\{|f(p) - f(q)|; \|[D, f]\| \leq 1\}.$$

Compare with the classical dual formula

$$d(p, q) = \text{Inf} \int_p^q (g_{\mu\nu} dx^\mu dx^\nu)^{1/2}$$

- ▶ Connes' notion of **noncommutative Riemannian manifold**, or spectral manifolds (A, \mathcal{H}, D) .

Discrete Dirac operators of Barrett-Glaser

- ▶ A discrete Dirac operator: $D : M_N(\mathbb{C}) \rightarrow M_N(\mathbb{C})$

$$D(X) = \{H, X\}, \quad H \in \mathcal{H}_N$$

Discrete Dirac operators of Barrett-Glaser

- ▶ A discrete Dirac operator: $D : M_N(\mathbb{C}) \rightarrow M_N(\mathbb{C})$

$$D(X) = \{H, X\}, \quad H \in \mathcal{H}_N$$

- ▶ Barrett-Glaser ([toy model for quantum gravity](#)): replace integration over metrics to integration over Dirac operators:

$$Z = \int_{\text{metrics}} e^{-S(g)} D(g) \Rightarrow \int_{\text{Diracs}} e^{-S(D)} dD$$

Discrete Dirac operators of Barrett-Glaser

- ▶ A discrete Dirac operator: $D : M_N(\mathbb{C}) \rightarrow M_N(\mathbb{C})$

$$D(X) = \{H, X\}, \quad H \in \mathcal{H}_N$$

- ▶ Barrett-Glaser ([toy model for quantum gravity](#)): replace integration over metrics to integration over Dirac operators:

$$Z = \int_{\text{metrics}} e^{-S(g)} D(g) \Rightarrow \int_{\text{Diracs}} e^{-S(D)} dD$$

- ▶ Consider a quartic action

$$S(D) = g \text{Tr}(D^2) + \text{Tr}(D^4)$$

Multitrace and multimatrix unitary invariant ensembles

- ▶ Computation shows:

$$S(D) = 2N(g \operatorname{Tr} H^2 + \operatorname{Tr} H^4) + 2g(\operatorname{Tr} H)^2 + 8 \operatorname{Tr} H \operatorname{Tr} H^3 + 6(\operatorname{Tr} H^2)^2$$

Barrett-Glaser studied these models via MCMC. What about their analytic treatment? This was my starting point.

Multitrace and multimatrix unitary invariant ensembles

- ▶ Computation shows:

$$S(D) = 2N(g \operatorname{Tr} H^2 + \operatorname{Tr} H^4) + 2g(\operatorname{Tr} H)^2 + 8 \operatorname{Tr} H \operatorname{Tr} H^3 + 6(\operatorname{Tr} H^2)^2$$

Barrett-Glaser studied these models via MCMC. What about their analytic treatment? This was my starting point.

- ▶ A unitary invariant function $F(H)$ on \mathcal{H}_N is of the form

$$F(H) = P(\operatorname{Tr}(H), \operatorname{Tr}(H^2), \dots, \operatorname{Tr}(H^N))$$

for a polynomial P in N -variables.

- ▶ Our models are all multitrace and mostly multimatrix, but unfortunately most results in RMT are for functions $F(H) = \operatorname{Tr}(V(H))$.

Genus expansion, summing over discrete surfaces/fat graphs ('t Hooft)

- ▶ Fix a polynomial $V(x) = \sum \frac{t_k}{k} x^k$. Consider the matrix integral

$$Z_N = \int_{\mathcal{H}_N} e^{-N\text{Tr}(V(H))} dH,$$

- ▶ **Topological expansion** of $F_N = \log Z_N$

$$F_N = \sum_{g \geq 0} (N)^{2-2g} F_g, \quad F_g = \sum_{[\mathcal{M}] \in \mathbb{M}_\emptyset^g} \text{weight}(\mathcal{M})$$

where \mathbb{M}_\emptyset^g = set of isomorphism classes of the Feynman weighted connected closed maps.

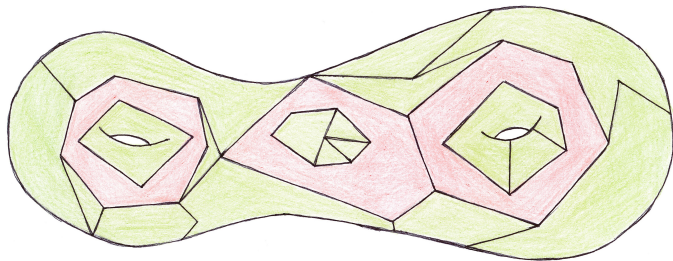


Figure: A polygonalization of a genus 2 surface

Genus expansion leads to a quick proof of the Wigner law, links with geometry of moduli spaces of curves, topological recursion (Eynard-Orantin), 2d gravity, recursion formula for volumes of moduli spaces of Riemann surfaces (Mirzakhani recursion).

Eynard-Orantin topological recursion (Tutte recursion)

$$\omega_{g,n} = K * \omega_{g-1,n+1} + \sum K * \omega_{g_1,n_1} \omega_{g_2,n_2}$$

Figure: Source: Wikipedia, by B. Eynard

- ▶ $\omega_{g,n}$ is a meromorphic symmetric n -form on Σ , recursion is on $-\xi(\Sigma) = 2g - 2 + n$.
- ▶ Example: Mirzakhani recursion for Weil-Petersson volumes of moduli spaces, $\omega_{0,1} = \frac{4}{\pi} z \sin(\pi z) dz$ and $\omega_{0,2}(z_1, z_2) = \frac{dz_1 dz_2}{(z_1 - z_2)^2}$

Resolvent technique

- ▶ Green function (Stieljes transform)

$$G(z) = \frac{1}{N} \langle \text{Tr} \frac{1}{z - H} \rangle = \int \rho_N(\lambda) \frac{d\lambda}{z - \lambda}$$

Is a holomorphic function on $\mathbb{C} \setminus \text{supp}(\rho_N)$.

- ▶ Can recover the eigenvalue probability density $\rho(x)$ from $G(x)$:

$$-2\pi i \rho_N(x) = \lim_{\varepsilon \rightarrow 0^+} [G(x + i\varepsilon) - G(x - i\varepsilon)]$$

- ▶ For single trace models with potential $V(x)$, It satisfies a differential equation

$$G^2(x) - V'(x)G(x) + \frac{1}{N}G'(x) + P(x) = 0.$$

- ▶ Can drop N for large N limit

$$G(x) = \frac{1}{2}(V'(x) - \sqrt{V'(x)^2 - 4P(x)})$$

(Since $G(x) \sim \frac{1}{x}$ for large x , negative sign is chosen).

- ▶ For $V(x) = \frac{1}{2}x^2$ this gives the semicircle law. For small enough variations support of $\rho(x)$ is an interval, but in general will be a union of intervals (multi cut regime).

The saddle point equation

- ▶ Our models are unitary invariant but not single trace, so a slightly different approach is needed.

$$Z = \int_{\mathcal{H}_N} e^{F(H)} dH,$$

$$Z = C_N \int_{\mathbb{R}^N} e^{-N \sum_{i=1}^N V(\lambda_i) - \sum_{i,j=1}^N U(\lambda_i, \lambda_j)} \prod_{1 \leq i < j \leq N} (\lambda_i - \lambda_j)^2 d\lambda_1 \dots d\lambda_N,$$

$$V(s) = 2gs^2 + 2s^4,$$

$$U(s, t) = 2gst + 8st^3 + 6s^2t^2,$$

Euler-Lagrange equations

- ▶ The measure $\mu = \rho(x)dx$ is referred to as the **equilibrium measure** and it is the Borel probability measure that **minimizes** the energy functional $I(\mu) =$

$$\int V(s)d\mu(s) + \int \int U(s, t)d\mu(s)d\mu(t) - \int \int \log|s-t|d\mu(s)d\mu(t).$$

Euler-Lagrange equations

- ▶ The measure $\mu = \rho(x)dx$ is referred to as the **equilibrium measure** and it is the Borel probability measure that **minimizes** the energy functional $I(\mu) =$

$$\int V(s)d\mu(s) + \int \int U(s, t)d\mu(s)d\mu(t) - \int \int \log|s-t|d\mu(s)d\mu(t).$$

- ▶ Euler-Lagrange equations (fully justified by results of Percy Deift) shows:

$$\text{P.V.} \int_{\text{supp}\rho} \frac{\rho(s)}{s-x} ds = 2gx + 4x^3 + gm_1 + 4m_3 + 6m_2x.$$

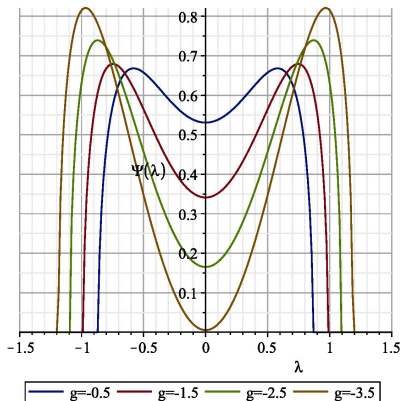


Figure: The equilibrium measure from the single cut analysis.

A precise critical value is found by setting $\rho(x) = 0$ at $x = 0$:

$$g_c = -\frac{5\sqrt{2}}{2} \approx -3.5,$$

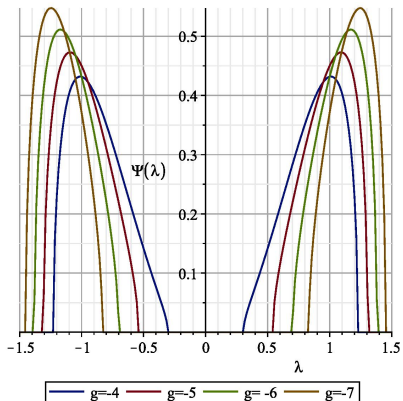


Figure: The equilibrium measure from the double cut analysis, $g < g_c$

References

- ▶ (M.K. and N. Pagliaroli) arXiv:2006.02891: Phase transitions in random noncommutative geometries.
- ▶ (S. Azarfar and M.K.): arXiv:1906.09362: Random finite noncommutative geometries and topological recursion.
- ▶ (M.K. and N. Pagliaroli): Spectral statistics of Dirac ensembles.
- ▶ (H. Hessam, M.K. and N. Pagliaroli): Bootstrapping random noncommutative geometries.